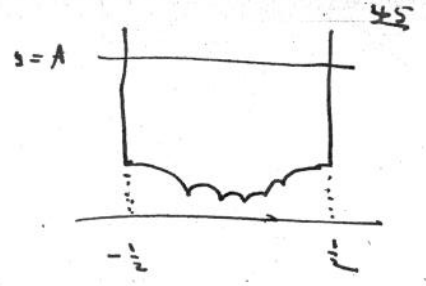


#16. Math 8, 1955

We had

$$E(z, s) = y^s + \varphi(s)y^{1-s} + \dots$$

$$\varphi(s) = \sqrt{\pi} \frac{\Gamma(s - \frac{1}{2})}{\Gamma(s)}, \quad \varphi(s)\varphi(1-s) = 1, \quad (\varphi(\frac{1}{2} + i\pi) = 1.$$



We define

$$\tilde{E}(z, s) = \begin{cases} E(z, s) & \text{for } y < A \\ E(z, s) - y^s - \varphi(s)y^{1-s} & \text{for } y \geq A \end{cases}$$

Let  $s \neq s', 1-s'$ . Consider

$$I = \iint_D \{ \tilde{E}(z, s) \Delta \tilde{E}(z, s') - \tilde{E}(z, s') \Delta \tilde{E}(z, s) \} dx dy.$$

Now  $\{y^2 \Delta + s(1-s)\} \tilde{E}(z, s) = 0$ . Hence we can replace  $\Delta \tilde{E}$  by  $-\frac{s(1-s)}{y^2} \tilde{E}$ . Hence

$$I = \{s(1-s) - s'(1-s')\} \iint_D \tilde{E}(z, s) \tilde{E}(z, s') \frac{dx dy}{y^2} \\ = (s'-s)(s+s'-1) \iint_D \tilde{E}(z, s) \tilde{E}(z, s') \frac{dx dy}{y^2}$$

By Green's formula, taking account of discontinuity along  $y=A$

$$I = \int_{\text{bdy.}} \{ \tilde{E}(z, s) \frac{\partial \tilde{E}(z, s')}{\partial n} - \tilde{E}(z, s') \frac{\partial \tilde{E}(z, s)}{\partial n} \} ds$$

Here  $y=A$  is part of boundary in opposite directions. Most of boundary parts cancel. Hence we will get

$$I = \left( \{y^s + y^{1-s}\varphi(s)\} \frac{d}{dy} \{y^{s'} + \varphi(s')y^{1-s'}\} - \{y^{s'} + \varphi(s')y^{1-s'}\} \frac{d}{dy} \{y^s + \varphi(s)y^{1-s}\} \right)_{y=A} \\ = \left( (s'-s) [y^{s+s'-1} - \varphi(s)\varphi(s')y^{1-s-s'}] \right. \\ \left. + (1-s-s') [\varphi(s')y^{s-s'} - \varphi(s)y^{s'-s}] \right)_{y=A}$$

$$I = (s'-s) \{ A^{s+s'-1} - \varphi(s)\varphi(s') A^{1-s-s'} \} + (s+s'-1) \{ \varphi(s) A^{s'-s} - \varphi(s') A^{s-s'} \}$$

$$= (s'-s)(s+s'-1) \iint_D \tilde{E}(z, s) \tilde{E}(z, s') \frac{dx dy}{y^2} \\ \frac{A^{s+s'-1} - \varphi(s)\varphi(s') A^{1-s-s'}}{s+s'-1} + \frac{\varphi(s) A^{s'-s} - \varphi(s') A^{s-s'}}{s'-s} = \iint_D \tilde{E}(z, s) \tilde{E}(z, s') \frac{dx dy}{y^2}$$

Let  $s = \sigma + ir$ ,  $s' = \sigma - ir$ ,  $r \neq 0$ ,  $\sigma \neq \frac{1}{2}$ . Then

$$\iint_{\mathcal{D}} |\tilde{E}(z, \sigma + ir)|^2 \frac{dx dy}{y^2} = \frac{A^{2\sigma-1} - |\varphi(\sigma + ir)|^2 A^{1-2\sigma}}{2\sigma-1} + \frac{\varphi(\sigma + ir) A^{2ir} - \overline{\varphi(\sigma + ir)} A^{-2ir}}{2ir}$$

With  $r \neq 0$ , let  $\sigma \rightarrow \frac{1}{2}$ . Since  $|\varphi(\frac{1}{2} + ir)| = 1$ ,  $\varphi$  is regular on  $\sigma = \frac{1}{2}$ .  
The first term has numerator

$$A^{2\sigma-1} - \varphi(\sigma + ir) \overline{\varphi(\sigma - ir)} A^{1-2\sigma}$$

with derivative

$$A^{2\sigma-1} \log A \cdot 2 - \varphi(\sigma + ir) \overline{\varphi(\sigma - ir)} A^{1-2\sigma} \log A (-2)$$

$$- A^{1-2\sigma} \{ \varphi(\sigma + ir) \varphi'(\sigma - ir) + \overline{\varphi(\sigma + ir)} \overline{\varphi'(\sigma - ir)} \}$$

$$\rightarrow 2 \log A + 2 \log A - \frac{\varphi(\frac{1}{2} + ir) \varphi'(\frac{1}{2} - ir) + \overline{\varphi(\frac{1}{2} + ir)} \overline{\varphi'(\frac{1}{2} - ir)}}{\varphi(\frac{1}{2} - ir) \varphi(\frac{1}{2} + ir)}$$

$$= 4 \log A - \left\{ \frac{\varphi'(\frac{1}{2} - ir)}{\varphi(\frac{1}{2} - ir)} + \frac{\overline{\varphi'(\frac{1}{2} + ir)}}{\overline{\varphi(\frac{1}{2} + ir)}} \right\}$$

$$= 4 \log A - 2 \Re \frac{\varphi'(\frac{1}{2} + ir)}{\varphi(\frac{1}{2} + ir)}$$

Hence the right side tends to

$$\frac{2}{\cancel{x}} \log A - \cancel{x} \Re \frac{\varphi'(\frac{1}{2} + ir)}{\varphi(\frac{1}{2} + ir)} + \frac{\varphi(\frac{1}{2} + ir) A^{2ir} - \overline{\varphi(\frac{1}{2} + ir)} A^{-2ir}}{2ir}$$

$$= \lim_{r \rightarrow \frac{1}{2}} \iint_{\mathcal{D}} |\tilde{E}(z, \sigma + ir)|^2 \frac{dx dy}{y^2}$$

Now, if  $E(z, s)$  has a pole  $s = \frac{1}{2} + ir$  with order  $\nu > 0$  then  $(s - \frac{1}{2} - ir)^\nu E(z, s)$  is regular at  $\frac{1}{2} + ir$  & this has the same limit  
 $\left( s - \frac{1}{2} - ir \right)^\nu \tilde{E}(z, s) \rightarrow u(z) \neq 0$  which would be an eigen-function  
 $\Rightarrow (y^2 \Delta + \frac{1}{4} + r^2) u = 0$ . But now  $\iint_{\mathcal{D}} |u|^2 \frac{dx dy}{y^2} = 0$   
 from the above since  $\nu > 0$ . Hence  $u = 0$  & this is a contradiction.  
 $\therefore$  No poles. Then we can replace the limit by  $|\tilde{E}(z, \frac{1}{2} + ir)|^2$ .  
 $\therefore$  No poles on  $\sigma = \frac{1}{2}$  except possibly  $s = \frac{1}{2}$ .

take  $\frac{1}{2} < \sigma \leq 1$  and let  $r \rightarrow 0$ . Suppose  $\varphi(z)$  is regular at  $\infty$ . Now differentiating again we get the limit

$$\frac{A^{2\sigma-1} - \varphi^2(\sigma) A^{1-2\sigma}}{2\sigma-1} + 2\varphi(\sigma) \log A - \varphi'(\sigma) = \lim_{r \rightarrow 0} \iint_D |\tilde{E}(z, \sigma + ir)|^2 \frac{dx dy}{y^2}$$

The left side  $\geq 0$ . With a pole of order  $\nu$  this can only happen if  $\nu < 2$   $\therefore \nu = 1$ . Then  $(s-\infty) \tilde{E}(z, s) \rightarrow u(z) \in L_2$

an eigen function. Except for the modular group  $\langle \begin{smallmatrix} a & b \\ c & d \end{smallmatrix} \rangle$  one can take  $A=1$  & this simplifies the above formula.

No poles for  $\frac{1}{2} < \sigma \leq 1$  except possibly a finite of simple ones and we can prove that  $s = \frac{1}{2}$  is also not a pole. One lets  $\sigma \rightarrow \frac{1}{2}$  in the above formula. To each pole at  $\sigma_0$  one gets an eigen function with  $\iint_D |u|^2 \frac{dx dy}{y^2} = c^2$

where  $(s-\sigma_0) \tilde{E}(z, s) \rightarrow c y^{1-\sigma_0} + \dots$

In the general case, one considers

$$\iint_D \tilde{E}_j(z, s, \chi) \overline{\tilde{E}_k(z, s', \chi)} \frac{dx dy}{y^2}$$

& cuts the fundamental region with several lines.

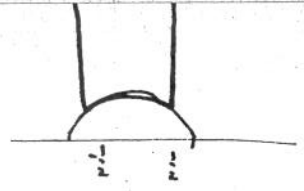
With eigen ~~functions~~  $E_1, E_2, \dots, E_n$

$$L(z, s, \chi) = \begin{pmatrix} E_1 \\ \vdots \\ E_n \end{pmatrix} = \Phi(s, \chi) E(z, 1-s, \chi)$$

One also uses  $A_1, \dots, A_n$  corresponding to  $\sigma_1, \dots, \sigma_n$ . One gets

$$\begin{aligned} & \iint_D \tilde{E}(z, \sigma + ir, \chi) \overline{\tilde{E}(z, \sigma + ir, \chi)} \frac{dx dy}{y^2} \\ &= \frac{1}{2\sigma-1} \left\{ A^{2\sigma-1} E - \varphi(\sigma + ir, \chi) \overline{\varphi(\sigma + ir, \chi)} A^{1-2\sigma} \right\} \\ & \quad + \frac{\varphi(\sigma + ir, \chi) A^{2ir} - \overline{\varphi(\sigma + ir, \chi)} A^{-2ir}}{2ir} \end{aligned}$$

Consider the modular group.



the Eisenstein series is

$$E(z, s) = \frac{1}{2} \sum_{(c,d)=1} \frac{y^s}{|cz+d|^{2s}} \quad \text{with the factor } \frac{1}{2} \text{ inserted so that we}$$

need no longer identify  $\frac{az+b}{cz+d}$ ,  $\frac{(-a)z+(-b)}{(-c)z+(-d)}$ . Now

$$2\zeta(2s) E(z, s) = \sum_{\substack{(c,d)=1 \\ n \geq 1}} \frac{y^s}{|cnz+dn|^{2s}} = \sum'_{m,n} \frac{y^s}{|mz+n|^{2s}} \\ = \sum' \frac{y^s}{\{(mx+ny)^2 + m^2y^2\}^s} \quad \text{Eisenstein } s\text{-function functional equation}$$

Pole at  $s=1$  only on right side. Singularities of  $E(z, s)$  correspond to zeros of  $\zeta(2s)$ .

$$E(z, s) = y^s + \Phi(s) y^{1-s} + \sum' e^{2\pi i x} \dots$$

$$\Phi(s) = \sqrt{\pi} \frac{\Gamma(s-\frac{1}{2})}{\Gamma(s)} \sum_{c>0} \frac{1}{c^{2s}} \varphi(c) \quad \text{Euler totient}$$

$$= \sqrt{\pi} \frac{\Gamma(s-\frac{1}{2})}{\Gamma(s)} \sum_{c>0} \frac{1}{c^{2s-1}} \sum_{d|c} \frac{\mu(d)}{d}$$

$$= \sqrt{\pi} \frac{\Gamma(s-\frac{1}{2})}{\Gamma(s)} \sum \sum \frac{\mu(d)}{d^{2s} m^{2s-1}}$$

$$= \sqrt{\pi} \frac{\Gamma(s-\frac{1}{2})}{\Gamma(s)} \frac{\zeta(2s-1)}{\zeta(2s)}$$

$$\frac{\pi^s}{\Gamma(s)} \sqrt{y} \int_0^\infty t^{s-\frac{1}{2}} e^{-\pi y(t+\frac{1}{2})} \frac{dt}{t} \sum_{c>0} \frac{1}{c^{2s}} \sum_{\substack{0 < d < c \\ (c,d)=1}} e^{2\pi i d/c} \mu(c)$$

$$= \frac{\pi^s}{\Gamma(s) \zeta(2s)} \sqrt{y} \int_0^\infty t^{s-\frac{1}{2}} e^{-\pi y(t+\frac{1}{2})} \frac{dt}{t}$$

We had  $\Phi(s)\Phi(1-s) = 1 \quad \therefore |\Phi(\frac{1}{2} + ir)| = 1 \quad \text{So}$

$$\left| \sqrt{\pi} \frac{\Gamma(ir)}{\Gamma(\frac{1}{2} + ir)} \frac{\zeta(2ir)}{\zeta(1+2ir)} \right| = 1$$

One can actually get the functional equation for  $\zeta(s)$  from  $\Phi(s)\Phi(1-s) = 1$  although it is not obvious.

In the more general case

$$\sum_{|c| \leq x} (x-|c|)^\alpha \sum_{0 \leq d \leq |c|} e^{2\pi i \frac{d}{x}} = \underbrace{\dots}_{\text{main term due to } \text{poles.}} + o\left(\frac{x^{\alpha+1}}{\log^\beta x}\right)$$

for certain  $\beta$

In the modular case

$$\sum_{n \leq x} (x-n)^\alpha \mu(n) = o\left(\frac{x^{\alpha+1}}{\log^\beta x}\right)$$

$$\sum_{n \leq x} \mu(n) = o\left(\frac{x}{\log^\beta x}\right)$$

Take the more general case with  $\chi(M) = 1$ . Now  $|\varphi(\frac{1}{2} + ir)| = 1$ . For  $\sigma > 1 + \epsilon$

$$\varphi(s) = \sqrt{\pi} \frac{\Gamma(s-\frac{1}{2})}{\Gamma(s)} \sum_{|c| > 0} \frac{1}{|c|^{2s}} \sum_{0 \leq d < |c|} 1 \quad \rightarrow \text{unif. bounded} \rightarrow 0 \text{ as } |s| \rightarrow \infty.$$

also

$$0 \leq \iint_{\mathcal{D}} |\tilde{E}(z, s)|^2 \frac{dx dy}{y^2} = \frac{A^{2\sigma-1} - |\varphi(\sigma+ir)|^2 A^{1-2\sigma}}{2\sigma-1} + \frac{\overline{\varphi(\sigma+ir)} A^{2ir} - \varphi(\sigma+ir) A^{-2ir}}{2ir}$$

taking  $A=1$

$$= \frac{1 - |\varphi(\sigma+ir)|^2}{2\sigma-1} + \frac{\overline{\varphi(\sigma+ir)} - \varphi(\sigma+ir)}{2ir}$$

let  $|r| \geq 1$ .

$$\leq \frac{1 - |\varphi(\sigma+ir)|^2}{2\sigma-1} + \frac{|\varphi(\sigma+ir)|}{|r|}$$

Hence, with  $\sigma > \frac{1}{2}$

$$1 - |\varphi|^2 + \frac{2\sigma-1}{r} |\varphi| \geq 0$$

$$|\varphi|^2 - \frac{2\sigma-1}{r} |\varphi| \leq 1$$

$$|\varphi| \leq \frac{\sigma-1}{|r|} + \sqrt{1 + \left(\frac{\sigma-1}{|r|}\right)^2} < 1 + \frac{2\sigma-1}{|r|}$$

$$|\varphi(\sigma+ir)| \leq 1 + \frac{2\sigma-1}{|r|}$$

$\therefore$  a bound for  $\varphi(s)$  for  $\Re s \geq \frac{1}{2}$ .

If there are real poles between  $\frac{1}{2}$  & 1 call them

$\frac{1}{2} < \sigma_1, \dots, \sigma_k \leq 1$ . Consider  $\varphi(s) \prod_{j=1}^k \frac{s - \sigma_j}{s - (1 - \sigma_j)} = \varphi^*(s)$ . Then

$|\varphi^*(\frac{1}{2} + i\sigma)| = 1$ . And we finally get  $|\varphi^*(s)| \leq 1$  for  $\Re(s) \geq \frac{1}{2}$   
 since  $\varphi^*(s)$  is regular for  $\Re(s) > \frac{1}{2}$  & bounded for  $\Re(s) > \frac{1}{2}$ .  
 Now  $z = \frac{s-1}{s}$  maps the half plane into the unit circle.



If the zeros of the transformed functions are  $\rho_1, \rho_2, \dots$   
 then  $\prod |\rho_n|$  and  $\sum [1 - |\rho_n|]$  converge. This is because

$$\frac{1}{2\pi} \int_0^{2\pi} \log |f(re^{i\theta})| d\theta \leq \frac{1}{2} \log \frac{1}{2\pi} \int_0^{2\pi} |f(re^{i\theta})|^2 d\theta$$

$\beta > 0$ , pole at  $\frac{1}{2} - \beta + i$

Jensen's formula. Hence if  $\frac{1}{2} + \beta + i\gamma$ , on the zero of  $\varphi(s)$   
 not on the real line then  $\sum \{1 - |\rho_n|^2\}$  also converges &

$$|\rho| = \left| \frac{\frac{1}{2} + \beta + i\gamma - 1}{\frac{1}{2} + \beta + i\gamma} \right| = \left| \frac{-\frac{1}{2} + \beta + i\gamma}{\frac{1}{2} + \beta + i\gamma} \right|$$

$$1 - |\rho|^2 = 1 - \frac{(\frac{1}{2} - \beta)^2 + \gamma^2}{(\frac{1}{2} + \beta)^2 + \gamma^2} = \frac{2\beta}{(\frac{1}{2} + \beta)^2 + \gamma^2}$$

Also  $\varphi(s) \neq 0$  for  $\Re(s) \geq \frac{1}{2}$  &  $\beta \leq \frac{1}{2}$ . Hence

~~Consider~~  $\sum \frac{\beta}{1 + \gamma^2}$  converges.  $\therefore$  if  $\gamma$  is large,  $\beta$  is close to 0. Consider

$$\prod \frac{s - \rho}{s - (1 - \bar{\rho})}$$

With  $s = \sigma + i\tau$ ,  $\rho = \frac{1}{2} + \beta + i\gamma$  then

$$\frac{s - \rho}{s - (1 - \bar{\rho})} = \frac{(\sigma - \frac{1}{2} - \beta) + i(\tau - \gamma)}{(\sigma - \frac{1}{2} + \beta) + i(\tau + \gamma)} = 1 - \frac{2\beta}{(\sigma - \frac{1}{2} + \beta) + i(\tau + \gamma)}$$

Arranging according to  $\rho$  &  $\bar{\rho}$  we would get an absolutely convergent product. Then

$$\varphi^*(s) = e^{g(s)} \prod \frac{s - \rho}{s - (1 - \bar{\rho})}$$

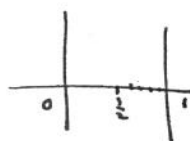
with  $g(s)$  reg. for  $\Re(s) > \frac{1}{2}$   
 &  $\Re g(s) = 0$  for  $\Re(s) = \frac{1}{2}$ .

With  $\varphi(s) = \frac{\text{entire function of order } \leq x}{\dots}$   $\Rightarrow g(s)$  is a polynomial of degree  $\leq 4$

Actually,  $g(s)$  is of the first degree,  $g(s) = a(s - \frac{1}{2})$  with a real  $a$ . This can be established by examining certain directions and comparing the size of the function. Actually  $a$  can be computed. Actually  $\sum_{|\rho| \leq T} \beta = O(T \log T)$  & this is best possible. In the <sup>general</sup> case, one uses the determinant of the matrix.

$$\varphi(s, \lambda) = \sqrt{\pi} \frac{\Gamma(s - \frac{1}{2})}{\Gamma(s)} L^{(0)}(s, \lambda)$$

$$\varphi(s, \lambda) \varphi(1-s, \lambda) = 1.$$



$$\frac{1}{2} < \sigma_1, \dots, \sigma_k \leq 1.$$

Drop  $\lambda$  in ~~dependent~~ argument. Let

$$\varphi^*(s) = \varphi(s) \prod_{j=1}^k \frac{s - \sigma_j}{s - 1 + \sigma_j}. \quad \text{Then}$$

$|\varphi^*(s)| = 1$  for  $s = \frac{1}{2} + i\tau$ . And  $\varphi^*(s)$  has no poles for  $\text{Re}(s) > \frac{1}{2}$ . Then

$|\varphi^*(s)| \leq 1$  for  $\text{Re}(s) \geq \frac{1}{2}$  on writing  $s = \sigma + i\tau$  all this we had

previously. Also  $\varphi^*(s) \varphi^*(1-s) = 1$ .

We had a function  $g(s)$ . Now things could be made so that

$g(s) + g(1-s) = 0$  or  $2\pi i$ . This means, since  $g(s)$  is of degree  $\leq 4$

$g(s) = \alpha(s - \frac{1}{2})^3 + \beta(s - \frac{1}{2}) + \gamma i\pi$ . Now  $\alpha, \beta$ , we deal by looking at what happens on  $\sigma = \frac{1}{2}$ . For large real  $s$ ,  $L^{(0)}(s, \lambda)$  is approx.

$$\lambda c_0^{-2s}$$

$$\varphi^*(s) \sim \sqrt{\pi} \frac{\lambda c_0^{-2s}}{\sqrt{s}}$$

Letting  $s = \frac{1}{2} + \sigma$  we have

$$\frac{s - \rho}{s - (1 - \bar{\rho})} = \frac{\frac{1}{2} + \sigma - \rho}{\frac{1}{2} + \sigma - 1 + \bar{\rho}}$$

$$= \frac{\sigma + \frac{1}{2} - (\frac{1}{2} + \beta + i\gamma)}{\sigma + \frac{1}{2} - 1 + (\frac{1}{2} + \beta + i\gamma)} = \frac{\sigma - \beta - i\gamma}{\sigma + \beta - i\gamma}$$

with the <sup>square</sup> absolute value

$$\frac{(\sigma - \beta)^2 + \gamma^2}{(\sigma + \beta)^2 + \gamma^2} = 1 - \frac{4\sigma\beta}{(\sigma + \beta)^2 + \gamma^2}$$

Consider  ~~$|\sigma| < \beta$~~ ,  ~~$|\sigma| \geq \beta$~~ .  $|\sigma| \leq T$ ,  $|\sigma| \geq T$ . Then

$$\prod_{|\sigma| < T} \left\{ 1 - \frac{4\sigma\beta}{(\sigma + \beta)^2 + \gamma^2} \right\} \geq \prod_{|\sigma| < T} \left\{ 1 - \frac{4\beta^2}{(\sigma + \beta)^2} \right\} \prod_{|\sigma| \geq T} \left\{ 1 - \frac{4\sigma\beta}{1 + \gamma^2} \right\}$$

$$\geq \prod_{|\sigma| < T} \left\{ 1 - \frac{\beta}{\sigma} \right\} \prod_{|\sigma| \geq T} \left( 1 - \frac{4\sigma\beta}{1 + \gamma^2} \right)$$

$$\geq \prod_{|\sigma| < T} e^{-2 \frac{\beta}{\sigma}} \prod_{|\sigma| \geq T} e^{-2 \frac{4\sigma\beta}{1 + \gamma^2}} \quad \text{larger}$$

$$\geq \prod_{|\sigma| < T} e^{-\frac{\beta}{\sigma}} \cdot e^{-8\sigma \sum_{|\sigma| \geq T} \frac{\beta}{1 + \gamma^2}} \rightarrow 0$$

$$\geq e^{-\frac{A}{T} T^{4+\frac{1}{2}}} \cdot e^{-8\sigma \sum_{1 \leq n \leq T} \frac{\beta}{n^2}} > e^{-o(\sigma)}$$

by suitable choice of  $T$ .

If  $\alpha \neq 0$  then  $g(s) \sim \alpha \sigma^3$  & we get a contradiction.

$\therefore \alpha = 0$ . Also we get  $\beta = -2 \log c_0$ . Hence

$$\varphi^*(s) = \pm c_0^{1-2s} \prod_{s-1+\rho} \frac{s-\rho}{s-1+\bar{\rho}}$$

$\rho$  are the zeros of  $\zeta(s)$  not on the real line.

From this it follows that the logarithm of  $\Pi$  is  $\sim \frac{1}{2} \log \sigma$ . Thus

$$\prod_{|s| \leq \sigma} \left\{ 1 - \frac{\alpha \beta}{(\sigma+s)^2 + \beta^2} \right\} > \frac{A}{\sigma}$$

$$\therefore \prod_{|s| \leq \sigma} \left( 1 - \frac{\beta}{\sigma} \right) > \frac{A}{\sigma}$$

$$\therefore \sum_{|s| \leq \sigma} \beta = O(\log \sigma)$$

Actually we can get  $A T \log T + B T + O(\log T)$ .

Consider

$$O \frac{\varphi^*(\frac{1}{2} + ir)}{\varphi^*(\frac{1}{2} + i\sigma)}$$

which arise from  $\iint_{\sigma} |\zeta(\frac{1}{2} + ir)|^2 \frac{dx dy}{y^2}$ .

Now

$$\frac{\varphi^*(s)}{\varphi^*(s)} = -2 \log c_0 + \sum_{\rho} \left( \frac{1}{s-\rho} - \frac{1}{s-1+\bar{\rho}} \right)$$

$$\frac{\varphi^*(\frac{1}{2} + ir)}{\varphi^*(\frac{1}{2} + i\sigma)} = -2 \log c_0 + \sum_{\rho} \left( \frac{1}{\frac{1}{2} + ir - \rho} - \frac{1}{\frac{1}{2} + ir - 1 + \bar{\rho}} \right) - \left( \frac{1}{\frac{1}{2} + i\sigma - \rho} - \frac{1}{\frac{1}{2} + i\sigma - 1 + \bar{\rho}} \right)$$

$$= -2 \log c_0 + \sum_{\rho} \left\{ \frac{1}{-\beta + i(r-\gamma)} - \frac{1}{\beta + i(r-\gamma)} \right\}$$

$$O \frac{\varphi^*(\frac{1}{2} + ir)}{\varphi^*(\frac{1}{2} + i\sigma)} = -2 \log c_0 + 2 \sum_{\rho} \frac{-\beta}{\beta^2 + (r-\gamma)^2}$$

Let

$$\omega(r) = -O \frac{\varphi^*(\frac{1}{2} + ir)}{\varphi^*(\frac{1}{2} + i\sigma)} = 2 \log c_0 + 2 \sum_{\rho} \frac{\beta}{\beta^2 + (r-\gamma)^2}$$

$$r-\gamma = \beta u$$

Now

$$\int_{-T}^T \omega(r) dr = 4T \log c_0 + 2 \sum \beta \int_{-T}^T \frac{dr}{\beta^2 + (r-\gamma)^2}$$

$$= \dots + 2 \sum \beta \int_{\frac{-T-\gamma}{\beta}}^{\frac{T-\gamma}{\beta}} \frac{\beta du}{\beta^2 + u^2 \beta^2}$$



$$\begin{aligned}
&= 4T \log c_0 + 2 \sum_{\substack{1 \leq \nu \leq T \\ \beta}} \int_{-\frac{T+\beta}{\beta}}^{\frac{T-\beta}{\beta}} \frac{du}{1+u^2} \\
&\leq 4T \log c_0 + 2 \sum_{\substack{1 \leq \nu \leq T \\ \beta}} \int_{-\infty}^{\infty} \frac{du}{1+u^2} + 2 \sum_{\substack{1 \leq \nu \leq T \\ \beta}} \frac{8T\beta}{\beta^2 + \beta^2} \\
&\leq 4T \log c_0 + 2 \cdot \pi \sum_{\substack{1 \leq \nu \leq T \\ \beta}} 1 + 2 \cdot 8T \sum_{\substack{1 \leq \nu \leq T \\ \beta}} \frac{\beta}{\beta^2 + \beta^2} \\
&= O(T) + 2\pi \sum_{\substack{1 \leq \nu \leq T \\ \beta}} 1 + O(T) \\
&= O(T) + O(T^{1+\epsilon}) = O(T^{1+\epsilon})
\end{aligned}$$

Similar things can be done for the determinant  $\det_{\lambda} \Phi(s, \chi)$  in the more general case. the matrix

$$\begin{aligned}
&\text{We had} \\
&\iint_{\mathcal{D}} \left| \tilde{E} \left( z, \frac{1}{2} + ir, \chi \right) \right|^2 \frac{dx dy}{y^2} = 2 \log A - R \frac{\chi'}{\chi} \left( \frac{1}{2} + ir, \chi \right) \\
&\quad + \frac{\varphi(\frac{1}{2} + ir, \chi) A^{2ir} - \varphi(\frac{1}{2} + ir, \chi) A^{-2ir}}{2ir}
\end{aligned}$$

The last term has ~~absolute~~ absolute value  $\leq \frac{1}{|r|}$ . From this we can estimate  $\tilde{E} \left( z, \frac{1}{2} + ir, \chi \right)$ . To do this we will need to investigate

$$H(z, z') = \frac{1}{4\pi} \int_{-\infty}^{\infty} e^{-it} |E(z, \frac{1}{2} + ir, \chi) \overline{E(z', \frac{1}{2} + ir, \chi)}| dt.$$

This is to be subtracted from  $K(z, z')$  the difference corresponding to the discrete spectrum.

#19. March 29, 1955

We had ~~for~~  $\frac{1}{2}$  for  $s = \sigma + ir$

$$\int_0^\infty |\tilde{E}(z, s, \lambda)|^2 \frac{dx dy}{y^2} = \frac{A^{2\sigma-1} |\varphi(\sigma+ir, \lambda)|^2 A^{1-2r}}{2\sigma-1} + \frac{\varphi(\sigma+ir, \lambda) A^{2ir} - \overline{\varphi(\sigma+ir, \lambda)} A^{-2ir}}{2ir}$$

$$= 2 \log A - 2R \frac{\varphi'}{\varphi} \left( \frac{1}{2} + ir, \lambda \right) + \dots$$

$$E(z, s, \lambda) = y^s + \varphi(s, \lambda) y^{1-s} + \sum' \alpha_n(s, \lambda) e^{2\pi n x} \sqrt{y} K_{s-\frac{1}{2}}(2\pi n y)$$

with  $K_{s-\frac{1}{2}}(y) = \int_0^\infty t^{s-\frac{1}{2}} e^{-y(t+\frac{1}{t})} \frac{dt}{t}$ .

Let  $s = \sigma + ir$ ,  $r \neq 0$ ,  $\frac{1}{2} \leq \sigma \leq \frac{3}{2}$ . Consider, Schwarz' inequality:

$$\left| \int_a^\infty f(y) K_{s-\frac{1}{2}}(y) dy \right| \leq \left\{ \int_a^\infty |f(y)|^2 dy \int_a^\infty |K_{s-\frac{1}{2}}(y)|^2 \frac{dy}{y} \right\}^{\frac{1}{2}}$$

We seek an  $f(y)$  with large  $\left| \int_a^\infty f(y) K_{s-\frac{1}{2}}(y) dy \right|$  + small  $\int_a^\infty |f(y)|^2 dy$ . Now

$$\begin{aligned} \int_a^\infty f(y) K_{s-\frac{1}{2}}(y) dy &= \int_a^\infty f(y) dy \int_0^\infty u^{s-\frac{1}{2}} e^{-y(u+\frac{1}{u})} \frac{du}{u} \quad u=yt \\ &= \int_a^\infty f(y) y^{s-\frac{1}{2}} dy \int_0^\infty t^{s-\frac{1}{2}} e^{-y^2 t - \frac{1}{t}} \frac{dt}{t} \\ &= \int_0^\infty t^{s-\frac{1}{2}} e^{-a^2 t - \frac{1}{t}} \frac{dt}{t} \int_a^\infty f(y) y^{s-\frac{1}{2}} e^{-(y^2-a^2)t} dy \end{aligned}$$

Let  $f(y) = y^{\frac{1}{2}-s} e^{ir \log(y^2-a^2)}$

Then we get

$$\begin{aligned} &\int_0^\infty t^{s-\frac{1}{2}} e^{-a^2 t - \frac{1}{t}} \frac{dt}{t} \int_a^\infty y \cdot e^{ir \log(y^2-a^2) - (y^2-a^2)t} dy \\ &= \frac{1}{2} \int_0^\infty t^{s-\frac{1}{2}} e^{-a^2 t - \frac{1}{t}} \frac{dt}{t} \int_0^\infty e^{ir \log u - tu} du \quad \frac{\Gamma(1+ir)}{t+ir} \\ &= \frac{1}{2} \Gamma(1+ir) \int_0^\infty t^{s-\frac{1}{2}} e^{-a^2 t - \frac{1}{t}} \frac{dt}{t} \\ &= \frac{1}{2} \Gamma(1+ir) \int_0^\infty t^{s-\frac{1}{2}} e^{-a(t+\frac{1}{t})} \frac{dt}{t} \end{aligned}$$

$$\begin{aligned} \frac{1}{2} K_{s-\frac{1}{2}}(a) &= \frac{1}{2} \int_0^\infty t^{s-\frac{1}{2}} e^{-a(t+\frac{1}{t})} \frac{dt}{t} \geq \int_1^\infty e^{-a(t+\frac{1}{t})} \frac{dt}{t} \geq \int_1^\infty e^{-a(t+1)} \frac{dt}{t} \\ &\geq e^{-a} \int_1^\infty t^{-1} e^{-at} dt = e^{-a} \int_a^\infty t^{-1} e^{-t} dt \\ &\geq e^{-a} \int_a^\infty \frac{1}{2} e^{-t} dt = e^{-a} \frac{1}{2} \int_a^\infty e^{-2t} dt = e^{-a} \frac{1}{2} e^{-2a} \end{aligned}$$

$$\left| \int_a^\infty f(y) K_{s-\frac{1}{2}}(y) dy \right| \geq \frac{1}{3} a^{\frac{1}{2}-\sigma} |\Gamma(1+ir)| e^{-\frac{1}{2}a}$$

$$\geq c_1 |\Gamma(1+ir)| e^{-3a}$$

Now  $s \Gamma(s) \Gamma(1-s) = \frac{\pi}{\sin \pi s}$   $s = ir$   $\Gamma(1+s) \Gamma(1-s) = \frac{\pi s}{\sin \pi s}$   $s = ir$

$$|\Gamma(1+ir)|^2 = \frac{\pi ir}{\sin \pi ir} = \frac{\pi ir}{\frac{e^{-\pi r} - e^{\pi r}}{2i}} = \frac{-2\pi r}{e^{-\pi r} - e^{\pi r}} = \frac{2\pi r}{e^{\pi r} - e^{-\pi r}}$$

$$|\Gamma(1+ir)| = \sqrt{\frac{2\pi r}{e^{\pi r} - e^{-\pi r}}} \approx \sqrt{2} e^{-\frac{\pi}{2}|r|}$$

$$\left| \int_a^\infty f(y) K_{s-\frac{1}{2}}(y) dy \right| \geq c_1 c_2 e^{-\frac{\pi}{2}|r|} e^{-3a} = c e^{-\frac{\pi}{2}|r| - 3a}$$

Now

$$\begin{aligned} |K_{s-\frac{1}{2}}(y)| &\leq \int_0^\infty t^{\sigma-\frac{1}{2}} e^{-y(t+\frac{1}{t})} \frac{dt}{t} \\ &\leq \int_0^\infty t e^{-y(t+\frac{1}{t})} \frac{dt}{t} = \int_0^\infty e^{-y(t+\frac{1}{t})} dt \\ &\leq \int_0^\infty (1+\frac{1}{t^2}) e^{-y(t+\frac{1}{t})} dt \\ &\leq 2 \int_0^\infty e^{-y(t+\frac{1}{t})} dt \leq 2 \int_0^\infty e^{-y t} dt = \frac{2e^{-y}}{y} \end{aligned}$$

and  $|f(y)| = y^{\frac{1}{2}-\sigma}$  so for  $R \geq 1$

$$\begin{aligned} \left| \int_{-R}^R f(y) K_{s-\frac{1}{2}}(y) dy \right| &\leq \int_R^\infty y^{\frac{1}{2}-\sigma} \cdot 2e^{-y} y^{-1} dy = 2 \int_R^\infty y^{\frac{1}{2}-\sigma} e^{-y} dy \\ &\leq 2 \int_R^\infty e^{-y} dy = 2e^{-R} \text{ Let} \end{aligned}$$

$R = 3a + \frac{\pi}{2}|r| + \omega$ , large  $\omega$ . Then

$$\left| \int_a^R f(y) K_{s-\frac{1}{2}}(y) dy \right| \geq \frac{1}{2} e^{-\frac{\pi}{2}|r| - 3a}$$

$$\begin{aligned} \frac{1}{2} e^{-\frac{\pi}{2}|r| - 3a} &\leq \int_a^R f(y) K_{s-\frac{1}{2}}(y) dy \leq \left\{ \int_a^R |f(y)|^2 y dy \int_a^R |K_{s-\frac{1}{2}}(y)|^2 \frac{dy}{y} \right\}^{\frac{1}{2}} \\ &\leq \left\{ \int_a^R y^{2\sigma} dy \int_a^\infty |K_{s-\frac{1}{2}}(y)|^2 \frac{dy}{y} \right\}^{\frac{1}{2}} \\ &\leq \frac{R^{2\sigma}}{2} \left\{ \int_a^\infty |K_{s-\frac{1}{2}}(y)|^2 \frac{dy}{y} \right\}^{\frac{1}{2}} \end{aligned}$$

Hence

$$\begin{aligned} \int_a^\infty |K_{s-\frac{1}{2}}(y)|^2 \frac{dy}{y} &\geq c' \left( \frac{e^{-\frac{\pi}{2}|r| - 3a}}{(3a + \frac{\pi}{2}|r| + \omega)^{2\sigma}} \right)^2 \geq c'' \left( \frac{e^{-\frac{\pi}{2}|r| - 3a}}{(\frac{\pi}{2}|r| + 3a + \omega)^2} \right)^2 \\ &= c'' \frac{1}{(\frac{\pi}{2}|r| + 3a + \omega)^{4\sigma}} \geq \frac{c''}{e^{2\sigma(\frac{\pi}{2}|r| + 3a + \omega)}} \geq e^{-(\frac{\pi}{2}|r| + 3a)(2\sigma)} \end{aligned}$$

then

$$|\alpha_n(s, \lambda)|^2 \int_0^\infty |K_{s-\frac{1}{2}}(\pi(n)y)|^2 \frac{dy}{y}$$

$$\geq |\alpha_n(s, \lambda)|^2 e^{-(\frac{\pi}{2}(1+3\pi(n)d))(2+\epsilon)}$$

We estimate  $K_{s-\frac{1}{2}}(y)$   
 change path of integration to other line. then



$$|K_{s-\frac{1}{2}}(y)| \leq \int_0^\infty e^{-(\frac{\pi}{2}-\epsilon)r} t^{\sigma-\frac{1}{2}} e^{-y(t+\frac{1}{2})\sin \epsilon} \frac{dt}{t}$$

$$\leq e^{-(\frac{\pi}{2}-\epsilon)r} K_{\sigma-\frac{1}{2}}(y \sin \epsilon)$$

$$\leq e^{-(\frac{\pi}{2}-\epsilon)r} \cdot \frac{2 e^{-y \sin \epsilon}}{y \sin \epsilon} \leq \frac{2 e^{-(\frac{\pi}{2}-\epsilon)r} e^{-y \frac{2}{\pi} \epsilon}}{y \frac{2}{\pi} \epsilon}$$

$$\leq \frac{\pi}{\epsilon y} e^{-\frac{\pi}{2} r - \frac{2}{\pi} \epsilon y}$$

$$\leq \frac{4}{\epsilon y} e^{-\frac{\pi}{2} r - \frac{1}{2} \epsilon y}$$

+ in general

$$|K_{s-\frac{1}{2}}(y)| \leq \frac{4}{\epsilon y} e^{-(\frac{\pi}{2}-\epsilon)r - \frac{1}{2} \epsilon y}$$

$\epsilon > 0$ .

Month April 5, 1955

We had for  $s = \sigma + ir$

$$\int_0^{\infty} |K^{\sim}(z, s, x)|^2 \frac{dx dy}{y^2}$$

$$= \begin{cases} 2 \log A - 2R \frac{\varphi'}{\varphi} \left(\frac{1}{2} + ir\right) + \frac{\varphi\left(\frac{1}{2} + ir\right) A^{2ir} - \varphi\left(\frac{1}{2} + ir\right) A^{-2ir}}{2ir} & \sigma = \frac{1}{2} \\ \frac{A^{2\sigma-1} - |\varphi(\sigma + ir)|^2 A^{1-2\sigma}}{2\sigma-1} + \frac{\varphi(\sigma + ir) A^{2ir} - \varphi(\sigma + ir) A^{-2ir}}{2ir} & \sigma > \frac{1}{2} \end{cases}$$

Now the last term is bounded for  $s$  away from the poles on the real axis. The first term for  $\sigma > \frac{1}{2}$  is

$$\frac{(A^{2\sigma-1} - A^{1-2\sigma}) + A^{1-2\sigma} (1 - |\varphi(\sigma + ir)|^2)}{2\sigma-1}$$

with the first part  $( ) / 2\sigma-1$  bounded for fixed  $A$  if  $\frac{1}{2} < \sigma \leq \frac{1}{2}$ .

Hence, since we can take  $A \geq 1$  everything depends on  $\frac{1 - |\varphi(\sigma + ir)|^2}{2\sigma-1}$ .

i.e. the expression is  $O(1) + O\left(\frac{1 - |\varphi(\sigma + ir)|^2}{2\sigma-1}\right)$ .

If we have no <sup>real</sup> poles then, with  $\rho = \beta + i\gamma$ ,  $\beta < \frac{1}{2}$

$$\varphi(s) = \pm c_0^{1-2s} \prod_{\rho} \frac{s-1+\bar{\rho}}{s-\rho}$$

$$|\varphi(\sigma + ir)|^2 = c_0^{2(1-2\sigma)} \prod_{\rho} \frac{(\sigma-1+\beta)^2 + (r-\gamma)^2}{(\sigma-\beta)^2 + (r-\gamma)^2}$$

$$= c_0^{-4(\sigma-\frac{1}{2})} \prod_{\rho} \left( 1 - \frac{(2\sigma-1)(1-\beta)}{(\sigma-\beta)^2 + (r-\gamma)^2} \right)$$

If  $0 \leq \nu_k \leq 1$  then  $1 - \prod \nu_k \leq \sum (1 - \nu_k)$ . Hence

$$1 - |\varphi(\sigma + ir)|^2 \leq 1 - c_0^{-4(\sigma-\frac{1}{2})} + (2\sigma-1) \sum_{\rho} \frac{1-2\beta}{(\sigma-\beta)^2 + (r-\gamma)^2}$$

$$\frac{1 - |\varphi(\sigma + ir)|^2}{2\sigma-1} \leq \frac{1 - c_0^{-2(2\sigma-1)}}{2\sigma-1} + \sum_{\rho} \frac{1-2\beta}{(\frac{1}{2}-\beta)^2 + (r-\gamma)^2}$$

$$= O(1) + \sum_{\rho} \frac{1-2\beta}{(\frac{1}{2}-\beta)^2 + (r-\gamma)^2}$$

$$= O(1) + 2R \frac{\varphi'}{\varphi} \left(\frac{1}{2} + ir\right)$$

If there are real poles then  $\varphi^*$  satisfies this inequality. and <sup>58</sup>

$$\frac{1-|\varphi|^2}{2\sigma-1} = \frac{1-|\varphi^*|^2}{2\sigma-1} + \frac{|\varphi^*|^2 - |\varphi|^2}{2\sigma-1}$$

The last term is uniformly bounded away from the real poles.

Hence the estimate for  $\iint_{\Omega} |\tilde{E}(z, s, \chi)|^2 \frac{dx dy}{y^2}$  is

~~essentially~~  $O(\omega(r))$  with  $\int_{-R}^R \omega(r) dr = O(R^{k+\epsilon})$

and  $\omega(r) = -\frac{2R^{k+\epsilon}(\frac{1}{2} + \epsilon r)}{y^*(\frac{1}{2} + \epsilon r)}$ ; this holds away from the real

poles.

By filling in with more fundamental regions we can show that

$$\iint_{\substack{y \geq \delta \\ |x| \leq \frac{1}{2}}} |\tilde{E}(z, s, \chi)|^2 \frac{dx dy}{y^2} = O(\omega(r))$$

$$\tilde{E}(z, s, \chi) = y^s + \varphi(s) y^{1-s} + \sum' \alpha_n(s) e^{2\pi i n x} \sqrt{y} K_{s-\frac{1}{2}}(\pi |n| y)$$

Taking as term of two sides

$$|\alpha_n(s)|^2 \iint_{\substack{y \geq \delta \\ |x| \leq \frac{1}{2}}} |K_{s-\frac{1}{2}}(\pi |n| y)|^2 \frac{dx dy}{y}$$

$$= |\alpha_n(s)|^2 \int_{\pi |n| \delta}^{\infty} |K_{s-\frac{1}{2}}(y)|^2 \frac{dy}{y} = O(\omega(r))$$

Hence

$$|\alpha_n(s)|^2 \cdot e^{-\left(\frac{\pi}{2}|n| + 3\pi|n|\delta\right)(z+\epsilon)} = O(\omega(r))$$

$$|\alpha_n(s)|^2 = O(\omega(r) e^{\left(\frac{\pi}{2}|n| + \frac{3\pi}{2}|n|\delta\right)(z+\epsilon)})$$

$$|\alpha_n(s)| = O\left(\sqrt{\omega(r)} e^{\left(\frac{\pi}{2}|n| + \frac{3\pi}{2}|n|\delta\right)(z+\epsilon)}\right)$$

Let  $E^*(z, s, \chi) = \sum' \alpha_n(s) e^{2\pi i n x} \sqrt{y} K_{s-\frac{1}{2}}(\pi |n| y)$ . Then

$$\begin{aligned} E^*(z, s, \chi) &= O \sum' \sqrt{\omega(r)} e^{\left(\frac{\pi}{2}|n| + \frac{3\pi}{2}|n|\delta\right)(z+\epsilon)} e^{\sqrt{y}} \\ &= O \frac{\sqrt{\omega(r)}}{\sqrt{y}} e^{(z+\epsilon)\chi |n|} \sum' e^{i\omega |n| \delta} \frac{1}{|n|} e^{-\frac{\pi}{2}|n| \delta y} \end{aligned}$$

$$= O \sqrt{\omega(r)} \frac{e^{(\epsilon+\epsilon')r}}{\sqrt{y}} \sum' e^{-\frac{\pi}{4} \frac{r^2}{y}} e^{-|n|(\frac{\pi}{4} \epsilon' y - 10\delta)}$$

and we want  $\frac{\pi}{4} \epsilon' y - 10\delta > 0$   $r \sim y \gg \frac{20\delta}{\pi \epsilon'}$ . We take

$$y \geq a + \text{then } \overset{\text{take } \delta}{10\delta} \leq \frac{\pi}{4} \epsilon' a \approx \frac{\pi}{4} \epsilon' y \text{ so that}$$

$$\frac{\pi}{4} \epsilon' y - 10\delta \geq \frac{\pi}{4} \epsilon' y - \frac{\pi}{4} \epsilon' y = \frac{\pi}{4} \epsilon' y$$

then get

$$E^+(z, s, \chi) = O \sqrt{\omega(r)} e^{(\epsilon+\epsilon')r} \frac{1}{\sqrt{y}} \sum' e^{-|n| \frac{\pi}{4} \epsilon' y}$$

$$= O \sqrt{\omega(r)} e^{(\epsilon+\epsilon')r} \frac{1}{\sqrt{y}} e^{-\frac{\pi}{4} \epsilon' y}$$

$$= O \left( \sqrt{\omega(r)} e^{2\epsilon r} \frac{1}{\sqrt{y}} e^{-\frac{\pi}{4} \epsilon' y} \right)$$

$$= O \left( \sqrt{\omega(r)} e^{2\epsilon r - \frac{\pi}{4} \epsilon' y} \right) \text{ for } y \geq a > 0.$$

We had some time ago

$$K(z, z') = \sqrt{yy'} g\left(\log \frac{z}{y}\right).$$

let

$$H(z, z') = \frac{1}{4\pi} \int_{-\infty}^{\infty} h(r) E(z, \frac{1}{2} + ir, \chi) E(z', \frac{1}{2} - ir, \chi) dr$$

Suppose  $h(r) = O(e^{-\delta|r|})$  for some  $\delta > 0$ .

then have to examine  $\int_{-\infty}^{\infty} \omega(r) e^{k\epsilon r} e^{-\delta|r|} dr$

then  $\int_{-\infty}^{\infty} \omega(r) e^{-\frac{\delta}{2}|r|} dr$ . Break up into

$$\sum_{k=0}^{2^{k+1}} \text{having an estimate } \sum e^{-\frac{\delta}{2} 2^k} (2^{k+1})^{k+\epsilon}$$

$$\text{since } \int_{-R}^R \omega(r) dr = O(R^{k+\epsilon}) \quad O \sum e^{-\frac{\delta}{2} 2^k} 2^{5k} \text{ we}$$

get convergence.

best time -

$h(r)$  even, regular for  $|h(r)| \leq \frac{1}{2} + \epsilon$

$$h(r) = O(e^{-\delta|r|})$$

$$H(z, z') = \frac{1}{4\pi} \int_{-\infty}^{\infty} h(r) E(z, \frac{1}{2} + ir, \lambda) E(z', \frac{1}{2} - ir, \lambda) dr$$

$$K(z, z') = \sum_{M \in \Gamma} \chi(M) k(z, Mz') \quad k \text{ derived from } h$$

Then in the case of  $1$  and  $M \in \Gamma$

$$K(z, z') - H(z, z') = O((yy')^\alpha) \text{ for some } \alpha < \frac{1}{2}$$

With  $\varphi(s, \lambda)$  having  $\sigma$  as the smallest pole  $> \frac{1}{2}$ , then we can take  $\alpha = 1 - \sigma$ . Kernels of the above type commute among themselves, the spectrum is discrete, the eigenfunctions is square-integrable. We write

$$E(z, s, \lambda) = y^s + \varphi(s, \lambda) y^{1-s} + E^*(z, s, \lambda)$$

then

$$E^*(z, s, \lambda) = O(\sqrt{w(r)} e^{2\epsilon|r| - \frac{\epsilon}{2}y})$$

$$w(r) \text{ even, } > 0, \int_0^R w(r) dr = O(R^{1+\epsilon}).$$

depend on  $A$

$$\iint_{\mathfrak{D}} |E^*(z, \frac{1}{2} + ir, \lambda)|^2 \frac{dx dy}{y^2}$$

$$= 2 \log A - 2R \frac{\varphi'}{\varphi}(\frac{1}{2} + ir, \lambda) + \frac{\varphi(\frac{1}{2} + ir, \lambda) A^{2ir} - \varphi(\frac{1}{2} - ir, \lambda) A^{-2ir}}{2ir}$$

Now

$$\iint_{\substack{y \geq A \\ |x| \leq \frac{1}{2}}} |E^*(z, \frac{1}{2} + ir, \lambda)|^2 \frac{dx dy}{y^2} = O(w(r) e^{2\epsilon|r|}) \int_A^{\infty} \int_{-\frac{1}{2}}^{\frac{1}{2}} \frac{e^{-\epsilon y}}{y^2} dx dy$$

$$= O(w(r) e^{2\epsilon|r|} \frac{e^{-\epsilon A}}{A^2})$$

$$\therefore \iint_{\substack{|y| < A \\ \mathfrak{D}}} |E^*(z, \frac{1}{2} + ir, \lambda)|^2 \frac{dx dy}{y^2} = 2 \log A - 2R \frac{\varphi'}{\varphi}(\frac{1}{2} + ir, \lambda) + \frac{\varphi(\frac{1}{2} - ir, \lambda) A^{2ir} - \varphi(\frac{1}{2} + ir, \lambda) A^{-2ir}}{2ir} + O(\frac{w(r) e^{2\epsilon|r|}}{A^2})$$



$$\begin{aligned}
 \iint_{S_A} H(z, z) \frac{dx dy}{y^2} &= \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) \int_{S_A} |E(z, \frac{1}{2} + ir, x)|^2 \frac{dx dy}{y^2} \\
 &= \frac{\log A}{2\pi} \int_{-\infty}^{\infty} h(r) dr - \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) R \frac{\phi'}{P} (\frac{1}{2} + ir, x) dr \\
 &\quad + \frac{1}{8\pi i} \int_{-\infty}^{\infty} \frac{h(r)}{r} \{ \varphi(\frac{1}{2} + ir, x) A^{2ir} - \varphi(\frac{1}{2} + ir) A^{-2ir} \} dr \\
 &\quad + O\left(\frac{e^{-\frac{\delta}{2}A}}{A^2} \int_{-\infty}^{\infty} |h(r)| w(r) e^{+\epsilon |r|} dr\right)
 \end{aligned}$$

Let  $\epsilon = \frac{\delta}{P}$  then integrand is  $O(e^{-\frac{\delta}{2}|r|} w(r))$  & we get an estimate of  $O(e^{-\frac{\delta}{2}A})$  for the last term. The next to the last integral is  $O(\frac{1}{\log A})$  by partial integration. Hence

$$\begin{aligned}
 \iint_{S_A} H(z, z) \frac{dx dy}{y^2} &= \frac{\log A}{2\pi} \int_{-\infty}^{\infty} h(r) dr - \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) R \frac{\phi'}{P} (\frac{1}{2} + ir, x) dr \\
 &\quad + O\left(\frac{1}{\log A}\right) + O\left(e^{-\frac{\delta}{2}A}\right) \\
 &= O\left(\frac{1}{\log A}\right) = o(1) \quad \text{as } A \rightarrow \infty.
 \end{aligned}$$

$$k(z, z+n) = k\left(\left(\frac{z-z+n}{y}\right)^2\right) = k\left(\frac{n^2}{y^2}\right)$$



$$\begin{aligned}
 \iint_{S_A} \sum_{n=-\infty}^{\infty} k(z, z+n) \frac{dx dy}{y^2} &= \int_0^A \sum_{-\infty}^{\infty} k\left(\frac{n^2}{y^2}\right) \frac{dx dy}{y^2} = \sum_{-\infty}^{\infty} \int_0^A k\left(\frac{n^2}{y^2}\right) \frac{dy}{y^2} \\
 &= \sum_{-\infty}^{\infty} \int_0^{\frac{A}{|n|}} k(u^2) \frac{-\frac{n}{u^2}}{n^2/u^2} du \quad \left( y = \frac{n}{u} \right) \\
 &= 2 \sum_{1 \leq |n| \leq A} \frac{1}{|n|} \int_{\frac{n}{A}}^{\infty} k(u^2) du \quad \left( \frac{n}{A} \leq u \right) \\
 &= 2 \int_{\frac{1}{A}}^{\infty} k(u^2) du \sum_{1 \leq |n| \leq An} \frac{1}{|n|} = 2 \int_0^{\infty} k(u^2) \{ \log Au + \gamma + O\left(\frac{1}{Au}\right) \} du \\
 &= 2 \int_0^{\infty} k(u^2) \log u du + 2(\log A + \gamma) \int_0^{\infty} k(u^2) du \\
 &\quad + O\left(\frac{1}{A} \int_0^{\infty} \frac{|k(u^2)|}{\sqrt{u}} du\right) = O\left(\frac{1}{A}\right)
 \end{aligned}$$

$k$  is bounded near 0  
 $k = O\left(\frac{1}{u^{\frac{1}{2} + \epsilon}\right)$

$$2 \int_0^\infty k(u^2) du = 2 \int_0^\infty k(t) \cdot \frac{1}{2} \frac{1}{\sqrt{t}} dt = \frac{1}{\sqrt{t}} \int_0^\infty \frac{k(t)}{\sqrt{t}} dt = g(0)$$

$$= \frac{1}{2i\pi} \int_{-\infty}^\infty r(r) dr.$$

Hence the logarithmic term in A drops out. Letting  $A \rightarrow \infty$  we write

$$\iint_{\mathcal{D}} \{K(z, z) - H(z, z)\} \frac{dx dy}{y^2} = \dots \quad (\text{non parabolic term})$$

$$+ \frac{1}{2\pi} \int_{-\infty}^\infty r(r) R \frac{p'}{p} \left(\frac{1}{2} + ir, \chi\right) dr$$

$$+ \frac{\gamma}{2\pi} \int_{-\infty}^\infty h(r) dr + \delta g(0)$$

$$+ 2 \int_0^\infty k(u^2) \log u du.$$

$u^2 = t$

$$I = 2 \int_0^\infty k(u^2) \log u du = 2 \int_0^\infty k(t) \log t \cdot \frac{1}{2} \frac{1}{\sqrt{t}} dt = \frac{1}{2} \int_0^\infty k(t) \frac{\log t}{\sqrt{t}} dt$$

NW

$$\int_w^\infty \frac{k(t)}{\sqrt{t-w}} dt = Q(w) \quad Q(e^u + e^{-u} - w) = g(u)$$

$$k(t) = -\frac{1}{\pi} \int_t^\infty \frac{dQ(w)}{\sqrt{w-t}}$$

so

$$I = -\frac{1}{2\pi} \int_0^\infty \frac{\log t}{\sqrt{t}} dt \int_t^\infty \frac{dQ(w)}{\sqrt{w-t}} \quad 0 \leq t \leq w < \infty$$

$$= -\frac{1}{2\pi} \int_0^\infty \frac{dQ(w)}{w} \int_0^w \frac{\log t}{\sqrt{t(w-t)}} dt \quad t = wu$$

$$= -\frac{1}{2\pi} \int_0^\infty dQ(w) \int_0^1 \frac{\log w + \log u}{\sqrt{u(1-u)}} \sqrt{w} du$$

$$\int_0^1 \frac{du}{\sqrt{u(1-u)}} = \int_0^1 u^{-1/2} (1-u)^{-1/2} du = B\left(\frac{1}{2}, \frac{1}{2}\right) = \frac{\Gamma(\frac{1}{2})\Gamma(\frac{1}{2})}{\Gamma(1)} = \pi$$

or  $u = \sin^2 \theta$  so  $\int_0^{\pi/2} \frac{2 \sin \theta \cos \theta d\theta}{\sin \theta \cos \theta} = \pi$ . Hence

$$I = -\frac{1}{2\pi} \int_0^\infty dQ(w) \left\{ \pi \log w + \int_0^1 \frac{\log u}{\sqrt{u(1-u)}} du \right\}$$

NW  $\int_0^1 u^{s-1} (1-u)^{\frac{1}{2}-1} du = \frac{\Gamma(s)\Gamma(\frac{1}{2})}{\Gamma(s+\frac{1}{2})} = \frac{\sqrt{\pi} \Gamma(s)}{\Gamma(s+\frac{1}{2})}$

Diff.

$$\int_0^1 u^{s-1} \log u \cdot u^{s-1} du = \sqrt{\pi} \frac{\Gamma(s+\frac{1}{2}) \Gamma'(s) - \Gamma(s) \Gamma'(s+\frac{1}{2})}{\Gamma^2(s+\frac{1}{2})}$$

$$\int_0^1 u^{-\frac{1}{2}} (1-u)^{-\frac{1}{2}} \log u du = \sqrt{\pi} \left\{ \frac{\Gamma'(\frac{1}{2}) - \sqrt{\pi} \Gamma'(1)}{\Gamma(\frac{1}{2})} \right\}$$

$$= \sqrt{\pi} \{ \Gamma'(\frac{1}{2}) + \sqrt{\pi} \gamma \}$$

$$= \pi \left\{ \frac{\Gamma'(\frac{1}{2})}{\Gamma(\frac{1}{2})} + \gamma \right\}$$

Now  $\Gamma(s) \Gamma(s+\frac{1}{2}) = A 2^{1-2s} \Gamma(2s)$

$$\Rightarrow \frac{\Gamma'(s)}{\Gamma(s)} + \frac{\Gamma'(s+\frac{1}{2})}{\Gamma(s+\frac{1}{2})} = -2 \log 2 + 2 \frac{\Gamma'(2s)}{\Gamma(2s)}$$

$$s = \frac{1}{2}$$

$$\frac{\Gamma'(\frac{1}{2})}{\Gamma(\frac{1}{2})} + \frac{\Gamma'(1)}{\Gamma(1)} = -2 \log 2 + 2 \frac{\Gamma'(1)}{\Gamma(1)}$$

$$\frac{\Gamma'(\frac{1}{2})}{\Gamma(\frac{1}{2})} = -2 \log 2 + \frac{\Gamma'(1)}{\Gamma(1)} = -\gamma - 2 \log 2$$

$$\int_0^1 \frac{\log u}{\sqrt{u(1-u)}} du = \pi(-\gamma - 2 \log 2 + \gamma) = -2\pi \log 2.$$

Here ~~But being give~~  $\pi(-2 \log 2)$ .

$$I = -\frac{1}{2\pi} \int_0^\infty (\pi \log w - 2\pi \log 2) dQ(w)$$

$$= -\log 2 \cdot g(0) = \frac{1}{2} \int_0^\infty \log w \cdot dQ(w)$$

$$w = e^u + e^{-u} - 2 = (e^{\frac{u}{2}} - e^{-\frac{u}{2}})^2$$

$$= -g(0) \log 2 - \int_0^\infty \log(e^{\frac{u}{2}} - e^{-\frac{u}{2}}) dg(u)$$

$$= -g(0) \log 2 - \int_0^\infty \log(e^{\frac{u}{2}} - e^{-\frac{u}{2}}) \cdot g'(u) du$$

$$= -g(0) \log 2 - \frac{1}{2} \int_0^\infty u g'(u) du - \int_0^\infty g'(u) \log(1 - e^{-u}) du$$

$$= -g(0) \log 2 - \frac{1}{2} \{ u g(u) - \int g(u) du \}_0^\infty - \int_0^\infty g'(u) \log(1 - e^{-u}) du$$

$$= -g(0) \log 2 + \frac{1}{2} \int_{-\infty}^\infty g(u) du - \int_0^\infty g'(u) \log(1 - e^{-u}) du$$

$$= -g(0) \log 2 + \frac{1}{2} h(0) - \int_0^\infty g'(u) \log(1 - e^{-u}) du$$

Now

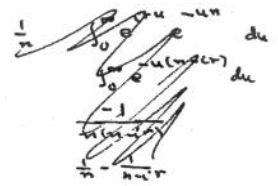
$$g(u) = \frac{1}{2\pi} \int_{-\infty}^\infty h(r) e^{iru} dr$$

$$\sim g'(u) = \frac{i}{2\pi} \int_{-\infty}^\infty h(r) e^{iru} r dr$$

$$\int_0^{\infty} g'(u) \log(1-e^{-u}) du = \frac{i}{2\pi} \int_0^{\infty} \log(1-e^{-u}) du \int_{-\infty}^{\infty} r h(r) e^{iru} dr$$

$$= \frac{i}{2\pi} \int_{-\infty}^{\infty} r h(r) dr \int_0^{\infty} e^{iru} \log(1-e^{-u}) du.$$

~~Since  $h(r)$  is even,  $r h(r)$  is odd~~



$$\int_0^{\infty} e^{iru} \log(1-e^{-u}) du = - \int_0^{\infty} e^{iru} \sum_{k=1}^{\infty} \frac{e^{-uk}}{k} du$$

$$= - \sum_{k=1}^{\infty} \frac{1}{k} \left\{ \frac{e^{u(ir-k)}}{ir-k} \right\}_0^{\infty} = \sum_{k=1}^{\infty} \frac{1}{k} \frac{1}{ir-k} = - \sum_{k=1}^{\infty} \frac{1}{k(k-ir)}$$

$$= \frac{1}{2} \left\{ \sum_{k=1}^{\infty} \frac{1}{k(k+ir)} - \sum_{k=1}^{\infty} \frac{1}{k(k-ir)} \right\} = \frac{1}{2} \left( \sum_{k=1}^{\infty} \frac{1}{k(k+ir)} + \sum_{k=1}^{\infty} \frac{1}{k(k-ir)} \right)$$

$$= \frac{1}{2} \sum_{k=1}^{\infty} \frac{-2ir}{k(k^2+r^2)} = \frac{1}{2} \left( \right)$$

$$= -ir \sum_{k=1}^{\infty} \frac{1}{k(k^2+r^2)} = \frac{1}{2} \left( \right)$$

$w_c$  had

$$K(z, z') = \sum_{M \in \Gamma} \chi(M) h(z, Mz')$$

only parabolic times

$$H(z, z') = \frac{1}{4\pi} \int_{-\infty}^{\infty} h(r) E(z, \frac{z}{2} + ir, t) \overline{E(z', \frac{z'}{2} + ir, t)} dr$$

Some things were not quite right in the earlier treatment so we go over it.

$$\iint_{\mathcal{D}} |\tilde{E}(z, \frac{z}{2} + ir, t)|^2 \frac{dx dy}{y^2} = 2 \log A - \frac{\varphi'}{\varphi}(\frac{z}{2} + ir, t) + \frac{\varphi(\frac{z}{2} - ir) A^{2ir} - \varphi(\frac{z}{2} + ir)}{2ir} + O(w(t) e^{\frac{2}{\delta} r - \frac{100}{\delta} A})$$

earlier had a factor of 2  
actually not so don't need it



because  $\varphi(1) \varphi(1/2) = 1$   
 $\frac{\varphi'(1)}{\varphi(1)} = \frac{\varphi'(1/2)}{\varphi(1/2)}$   
 $\frac{\varphi'(\frac{z}{2} + ir)}{\varphi(\frac{z}{2} + ir)} = \frac{\varphi'(\frac{z}{2} - ir)}{\varphi(\frac{z}{2} - ir)} = \left\{ \frac{\varphi'(\frac{z}{2} + ir)}{\varphi(\frac{z}{2} + ir)} \right\}$

$$\iint_{\mathcal{D}_A} H(z, z) \frac{dx dy}{y^2} = \frac{1}{2\pi} \log A \int_{-\infty}^{\infty} h(r) \frac{dr}{r} - \frac{1}{4\pi} \int_{-\infty}^{\infty} h(r) \frac{\varphi'}{\varphi}(\frac{z}{2} + ir) dr + \frac{1}{4\pi} \int_{-\infty}^{\infty} h(r) \frac{\varphi(\frac{z}{2} + ir) A^{2ir} - \varphi(\frac{z}{2} - ir) A^{-2ir}}{2ir} dr + O(e^{-\frac{5}{\delta} A})$$

The last integral  $J$  is O.K. except near  $r=0$  &  $\int_{-\infty}^{-\epsilon} + \int_{\epsilon}^{\infty} \rightarrow 0$

as  $A \rightarrow \infty$ .  $\therefore$  value depends on  $h(0) \tau \varphi(\frac{z}{2})$ . Consider then

$$I_{\epsilon} = \frac{1}{4\pi} \int_{-\infty}^{\infty} h(r) \varphi(\frac{z}{2}) e^{-\epsilon r^2} \frac{A^{2ir} - A^{-2ir}}{2ir} dr.$$

$$J - I_{\epsilon} = \frac{1}{4\pi} \int_{-\infty}^{\infty} \left\{ \frac{h(r) \varphi(\frac{z}{2} + ir) - h(r) \varphi(\frac{z}{2}) e^{-\epsilon r^2}}{2ir} A^{2ir} + \dots \right\} dr$$

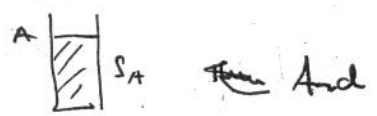
$\rightarrow 0$  as  $A \rightarrow \infty$  by the ~~Cauchy~~ Riemann-Lebesgue theorem.

And, with conditional convergence

$$I_0 = \lim_{\epsilon \rightarrow 0} \frac{1}{4\pi} \int_{-\infty}^{\infty} \frac{\sin(2r \log A)}{r} dr = \int_{-\infty}^{\infty} \frac{\sin u}{u} du = \pi$$

$$\therefore J \rightarrow I_0 = \pi h(0) \frac{\varphi(\frac{z}{2})}{4\pi} = \frac{1}{4} h(0) \varphi(\frac{z}{2}), \quad \varphi(\frac{z}{2}) = \pm 1.$$

$$\iint_{D_A} H(z, z) \frac{dx dy}{y^2} = g(z) \log A + \frac{1}{\pi} \varphi\left(\frac{z}{A}\right) h(z) - \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) \frac{\varphi'}{\varphi}\left(\frac{z}{z+ir}\right) dr + o(1) \quad \text{as } A \rightarrow \infty$$



$$\iint_{S_A} \sum' h(z, z+in) \frac{dx dy}{y^2} = (\log A + \gamma - \log z) g(z) + \frac{1}{\pi} h(z) - \int_0^{\infty} g'(u) \log(1-e^{-u}) du + o(1) \quad \text{as } A \rightarrow \infty$$

↑ Euler constant

Here

$$g'(u) = \frac{1}{2\pi} \int_{-\infty}^{\infty} ir h(r) e^{iru} dr$$

$$\begin{aligned} \mathcal{L} \left[ \sum' h(z, z+in) \frac{dx dy}{y^2} \right] &= - \int_0^{\infty} g'(u) \log(1-e^{-u}) du = - \frac{1}{2\pi} \int_{-\infty}^{\infty} ir h(r) dr \int_0^{\infty} e^{-ru} \log(1-e^{-u}) du \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} ir h(r) \sum_{k=1}^{\infty} \frac{1}{k} \frac{1}{k-ir} dr \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) \sum_{k=1}^{\infty} \frac{ir}{k(k-ir)} dr \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) \sum_{k=1}^{\infty} \left( \frac{1}{k-ir} - \frac{1}{k} \right) dr \end{aligned}$$

Now

$$\frac{1}{\Gamma(z)} = e^{\gamma z} \prod_{k=1}^{\infty} \left(1 + \frac{z}{k}\right) e^{-\frac{z}{k}}$$

↑ Euler constant

$$\begin{aligned} - \frac{\Gamma'(z)}{\Gamma(z)} &= \frac{1}{z} + \gamma + \sum_{k=1}^{\infty} \left\{ \frac{1}{1+\frac{z}{k}} - \frac{1}{k} \right\} \\ &= \frac{1}{z} + \gamma + \sum_{k=1}^{\infty} \left( \frac{1}{k+z} - \frac{1}{k} \right) \end{aligned}$$

$$\begin{aligned} \Gamma(z+1) &= z\Gamma(z) \\ \frac{\Gamma'(z+1)}{\Gamma(z+1)} &= \frac{1}{z} + \frac{\Gamma'(z)}{\Gamma(z)} \end{aligned}$$

$$\text{So } \sum_{k=1}^{\infty} \left( \frac{1}{k+z} - \frac{1}{k} \right) = -\gamma - \frac{1}{z} - \frac{\Gamma'(z)}{\Gamma(z)} = -\gamma - \frac{\Gamma'(z+1)}{\Gamma(z+1)}$$

$$\mathcal{L} \left[ \sum' h(z, z+in) \frac{dx dy}{y^2} \right] = \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) \left\{ -\gamma - \frac{\Gamma'(1+ir)}{\Gamma(1+ir)} \right\} dr$$

by evenness of  $h(r)$  we get

$$\mathcal{L} \left[ \sum' h(z, z+in) \frac{dx dy}{y^2} \right] = -\gamma g(z) - \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) \frac{\Gamma'}{\Gamma}(1+ir) dr$$

$$\therefore \int_{\Sigma_A} \Sigma^{-1} h(z, z+n) \frac{dx dy}{y^2} = (\log A - \log 2) g(0) + \frac{1}{4} h(0) - \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) \frac{\Gamma'(1+ir)}{\Gamma(1+ir)} dr + o(1)$$

$$\therefore \iint_D (K-H) \frac{dx dy}{4y^2} = \frac{1}{4\pi} \int_{-\infty}^{\infty} h(r) \frac{\varphi'}{\varphi} \left(\frac{1}{2} + ir\right) dr - \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) \frac{\Gamma'}{\Gamma}(1+ir) dr + \frac{1}{4} (1 - \varphi(\frac{1}{2})) h(0) - g(0) \log 2$$

on letting  $A \rightarrow \infty$ .

Letting the eigenvalues of the discrete spectrum be  $\frac{1}{2} + i\gamma$  (even allowing  $\gamma$  to be purely imaginary). By the trace formula, without the factor 2

$$\sum_{\gamma} h(\gamma) = \frac{1}{2} \{1 - \varphi(\frac{1}{2})\} h(0) + \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) \frac{-\varphi'}{\varphi} \left(\frac{1}{2} + ir\right) dr$$

$$= \frac{1}{2} \int_{-\infty}^{\infty} r \frac{e^{\pi r} - e^{-\pi r}}{e^{\pi r} + e^{-\pi r}} h(r) dr - \frac{1}{\pi} \int_{-\infty}^{\infty} h(r) \frac{\Gamma'}{\Gamma}(1+ir) dr$$

(elliptic)

$$+ \dots = g(0) \log 2 + 2 \sum_{\{P\}} \sum_{k=1}^{\infty} \frac{\chi^k \{P\} \log N\{P\}}{(N\{P\})^k - (N\{P\})^{-k}} \left. \vphantom{\sum_{\{P\}}} \right\} g(k \log N\{P\})$$

In the general case have the matrix  $\Phi(s, \chi)$  & logarithmic deriv. of determinant of  $\Phi(s, \chi)$  in place of  $\frac{\varphi'}{\varphi}(\frac{1}{2} + ir)$ .  
 Certain ~~factor~~ <sup>term</sup> hold up a factor of  $x$  in the general case.

~~This holds for~~  ~~$\frac{\varphi'}{\varphi}$~~   ~~$\frac{\Gamma'}{\Gamma}$~~  This holds for  $h(r)$  even  $\frac{\varphi'}{\varphi}$   $h(r) = O(e^{-\delta|r|})$

for  $|h(r)| \leq \frac{1}{2} + \epsilon$ . We want to extend this to more general  $h(r)$  & for this need <sup>an estimate for</sup> the number of zeros. This depends on the new integral  $\int_{-\infty}^{\infty} h(r) \frac{-\varphi'}{\varphi}(\frac{1}{2} + ir) dr$ . Now

$$\varphi(s) = C_0^{-1-2s} \prod_p \frac{s-1+p}{s-p}$$

$\rho$  a pole,  $\rho = \frac{1}{2} - \beta + i\gamma$  with  $\beta \geq 0$  except for a finite #

$$\frac{\varphi'(s)}{\varphi(s)} = -2 \log C_0 + \sum_p \left\{ \frac{1}{s-1+p} - \frac{1}{s-p} \right\}. \text{ Then}$$

$$\frac{\varphi'}{\varphi}(\frac{1}{2}+ir) = -2 \log C_0 + \sum \left\{ \frac{1}{-\beta+i(r-\lambda)} - \frac{1}{\beta+i(r-\lambda)} \right\}$$

$$= -2 \log C_0 + \sum \frac{-\beta-i(r-\lambda) - \beta+i(r-\lambda)}{\beta^2+(r-\lambda)^2}$$

$$= -2 \log C_0 + \sum \frac{-2\beta}{\beta^2+(r-\lambda)^2}$$

$C_0 \geq 1$

$$-\frac{\varphi'}{\varphi}(\frac{1}{2}+ir) = 2 \log C_0 + 2 \sum \frac{\beta}{\beta^2+(r-\lambda)^2}$$

let  $h(r) = e^{-r^2/R^2}$  with  $R$  large. then  $\int_{-\infty}^{\infty} h(r) dr = O(R^2)$

The result is that  $\sum_{|r| \leq R} h(r) \sim \frac{1}{2\pi} \int_{-\infty}^{\infty} h(r) \frac{-\varphi'}{\varphi} dr = O(R^2)$

$$\therefore \sum_{|r| \leq R} h(r) \geq \sum_{|r| \leq R} \frac{1}{e} = \frac{1}{e} N(R). \text{ So}$$

$$\frac{1}{e} N(R) + \frac{1}{e} \int_{-R}^R \frac{-\varphi'}{\varphi}(\frac{1}{2}+ir) dr = O(R^2)$$

$$\therefore N(R) = O(R^2) + \int_{-R}^R \frac{-\varphi'}{\varphi}(\frac{1}{2}+ir) dr = O(R^2).$$

So  $h(r) = O(\frac{1}{r^2+\epsilon})$  the right side of the trace formula is meaningful. Now use in place of such an  $h(r)$  the function  $h_\epsilon(r) = h(r) e^{-\epsilon r^2}$  & on letting  $\epsilon \rightarrow 0$  get the result

general  
for  $h(r) = O(\frac{1}{r^2+\epsilon})$

We had 
$$Z_n^*(s, \chi) = \prod_{\{P\}} \prod_{v=0}^{\infty} (1 - \chi(P)^v (N(P))^{-s-v})$$

~~we took~~ 
$$h(r) = \frac{1}{(s-\frac{1}{2})^2+r^2} - \frac{1}{(a-\frac{1}{2})^2+r^2}, \quad a > 1, \quad \Re(s) > 1$$

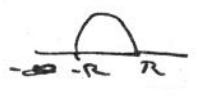
$$= O(\frac{1}{r^4})$$

the new term  $\int_{-\infty}^{\infty} h(r) \frac{-\varphi'}{\varphi} dr$



gives a <sup>new</sup> contribution. This depends on

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{(s-\frac{1}{2})^2 + r^2} \frac{2\rho}{\beta^2 + (r-\lambda)^2} dr$$



which, by contour integration, can be evaluated. We set

$$\frac{1}{2} \sum_{\rho} \left( \frac{1}{s-\frac{1}{2}} \left\{ \frac{1}{s-\rho} + \frac{1}{s-\bar{\rho}} \right\} - \frac{1}{a-\frac{1}{2}} \left\{ \frac{1}{a-\rho} + \frac{1}{a-\bar{\rho}} \right\} \right)$$

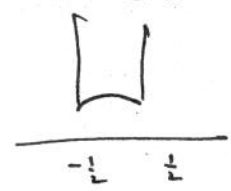
The terms  $2 \sum_{\{P\}} \sum_{k=1}^{\infty} \dots$  gave rise to  $\frac{1}{s-\frac{1}{2}} \frac{Z'}{Z}(s, \gamma)$ .

The zeros of  $\varphi(s)$  are not symmetric about  $\frac{1}{2}$  so there is not a real symmetry for  $Z(s)$ . But it turns out that

$$\frac{Z(1-s)}{Z(s)} = \varphi(s) \Psi(s) \quad \text{where } \Psi(s) \text{ has zeros at integers and at half integers.}$$

For the modular group we use  $\frac{Z(s)}{S(s)}$  which has a nice functional equation.

Consider the full modular group  $\Gamma$  ~~is that~~, let  $M = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma$   
 with  $f(Mz) = f(z)$  invariant under  $\Gamma$ . let  
 $f^*(z) = \begin{cases} f(z) & \text{for } z \in \mathcal{D} \\ 0 & \text{for } z \notin \mathcal{D} \end{cases}$



Then  $f(z) = \sum_{M \in \Gamma} f^*(Mz)$ .

Now let  $\begin{vmatrix} a & b \\ c & d \end{vmatrix} = n > 0$  &  $M_n = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$ . Consider

$\sum_{M_n} f^*(M_n z) = T_n f(z)$

which is automorphic. These are actually the Hecke operators. They commute &  $T_m T_n = T_{mn}$  for  $(m, n) = 1$ .

We had  $K(z, z') = \sum_{M \in \Gamma} k(z, Mz')$ . Now we use  $\sum_{M_n} k(z, M_n z')$ .

& get commuting operators, symmetric ones, etc.

unitary inner product.  $F(Mz) = \chi(M) F(z)$ . Equivalently, we could have dealt with  $(?)$   
 $\varphi(z, \mathbb{K}) = \varphi(Mz, \chi(M) \mathbb{K})$  with  $\mathbb{K}$  with  $\mathbb{K}$  an  $r$ -dimensional space.  
 & the diff. equation  $y^2 \Delta \varphi + (\frac{1}{2} + r^2) \varphi = 0$ .

Automorphic forms satisfy  $f(z) = \begin{pmatrix} \chi(M) \\ (cz+d) \end{pmatrix}^{-k} f(Mz)$ ;  $k$  is an integer although it might only be restricted to real values. If  $z' = Mz$  then  $y' = \frac{y}{(cz+d)^2} = \frac{y}{(c\bar{z}+d)(c\bar{z}'+d)}$  &  $\therefore (c\bar{z}+d)(c\bar{z}'+d) y' = y$ .

Then letting  $\mu(z) = y^{\frac{k}{2}} f(z)$  we have

$\mu(z) = \chi(M) \left( \frac{c\bar{z}'+d}{c\bar{z}+d} \right)^{\frac{k}{2}} \mu(Mz)$   
 ↑ absolute value 1.

For  $G$ :  $z \rightarrow mz$  &  $t \rightarrow t - \arg(cz+d) + \alpha$ . For  $\mu$ :  $z \rightarrow -\bar{z}$  &  $t \rightarrow -t$ . universal covering group

We can establish that with  $b = (s, t)$  we can make  $b = (s', t')$   
 $\left\{ \begin{matrix} b \rightarrow \mu b \\ b \rightarrow \nu b \end{matrix} \right\}$  with same  $\mu$ . As an invariant we have  $\frac{|t-z'|^2}{|t'|^2}$  as before.  
 Now  $\arg(z-\bar{z}') \rightarrow \arg\left(\frac{az+tb}{cz+d} - \frac{a\bar{z}'+t}{c\bar{z}'+d}\right) = \arg \frac{z-\bar{z}'}{(cz+d)(c\bar{z}'+d)}$

day

$$= \arg(z - \bar{z}') - \arg(cz + d) + \arg(cz' + d).$$

And  $t - t' \rightarrow t - t' - \arg(cz + d) + \arg(cz' + d)$ . Hence another invariant is  $t - t' - \arg(z - \bar{z}')$ . Under  $\mu$  this last changes sign. Hence we should deal with

$$k \left( \frac{|z - z'|^2}{y y'}, t - t' - \arg(z - \bar{z}') \right).$$

Then  $\frac{dx^2 + dy^2}{y^2}$  is an invariant form + so is  $dt - \frac{dx}{2y}$ .

This will not give a symmetric space in the sense of Cartan or using  $\frac{dx^2 + dy^2}{y^2} + k \left( dt - \frac{dx}{2y} \right)^2$  with  $k > 0$ . One gets  $\frac{\partial}{\partial t}$  & some thing else as invariant operators.

The eigenfunction would have to look like  $\varphi(z, t) = f(z) e^{i\alpha t}$

$\therefore \chi(M) = 1$  &  $k$  is an integer would let  $\varphi(z, t) = e^{i\alpha t} f(z)$ . From  $k(z, z')$  one forms  $K(z, z')$  & the rest of the theory. Nat much more difficult than before.

Consider, with  $\chi(M) = 1$ ,

$$f(z, \bar{z}) = (cz + d)^{-k} f\left(\frac{az + b}{cz + d}, \frac{a\bar{z} + \bar{b}}{c\bar{z} + d}\right)$$

$$\frac{\partial f}{\partial \bar{z}}(z, \bar{z}) = (cz + d)^{-k} (c\bar{z} + d)^{-2} \frac{\partial f}{\partial \bar{z}}(Mz, M\bar{z}).$$

$$\therefore \int_y \frac{\partial f}{\partial \bar{z}} = f^*(z)$$

then  $f^*(z) = (cz + d)^{2-k} f^*(Mz)$

Can do a similar thing to get  $(c\bar{z} + d)^{k-2}$  & only  $k$ 's between  $-1$  &  $1$  need be looked at. & by conjugation need only look at  $0 \dots 1$ .

Let  $K(z, z') = \sum \frac{1}{(cz + d)^k (Mz - \bar{z}')^k}$  & consider

$$f(z) = \lambda \int_{\mathcal{D}} \int_y \frac{1}{y^k} K(z, z') f(z') \frac{dx' dy'}{y'^2}$$

Can also use the Hecke ~~form~~ operators by using  $M_n$  in place of  $M$ . One could get the Ramanujan  $\tau(n)$  but not a better estimate than  $O(n^6)$ .

Part 5

Positive definite symmetric matrices  $Y > 0$ ;  $AYA' = G$   
 $|y| = 1$ .

Let  $y \rightarrow y^{-1}$  Take  $3 \times 3$ . Two fundamental operators.  
with  $|A| = 1$  &  $x$  a vector

$$\sum_{x'z=0} \frac{1}{(x' y x)^2 (z' y^{-1} z)^2}$$
 are eigen-functions, <sup>new type of</sup> zeta function

~~with~~ Dirichlet series of 2 variables. <sup>three</sup> Functional equations.

Also a discrete spectrum. One considers

$$k(\sigma(AyA'y_i^{-1}), \sigma(AyA'y_i^{-1})^2).$$

$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$   
 $\gamma = \frac{1}{\sqrt{1 - \beta^2}}$   
 $\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$   
 $\gamma = \frac{1}{\sqrt{1 - \beta^2}}$   
 $\gamma = \frac{1}{\sqrt{1 - \beta^2}}$

$$\left( \frac{1}{\sqrt{1 - \beta^2}} \right) = \left( \frac{1}{\sqrt{1 - \beta^2}} \right) \gamma$$

$$N(\alpha) = \gamma N(\beta)$$

$$u^2 = \alpha v^2 = \gamma^2 x^2 + \gamma^2 z^2$$

(72) reverse side